

A Direct Determination of the Crystal Structure of 2,3,4,6-Tetranitroaniline

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The symbolic addition procedure has been used to determine the crystal structure of 2,3,4,6-tetranitroaniline, $C_6H_3N_5O_8$. The crystals are monoclinic in space group $P2_1/c$ with $a_0 = 7.27$, $b_0 = 11.06$, $c_0 = 12.27 \text{ \AA}$ and $\beta = 98^\circ 47'$. The calculated density is 1.86 g.cm^{-3} with 4 molecules in the unit cell. The measured density is 1.87 g.cm^{-3} . The structure was refined, by both Fourier and full-matrix least-squares techniques, to a final R value of 0.08. The carbon–carbon bonds in the benzene ring are not 1.40 \AA . Those at the amine group carbon average 1.43 \AA and the remainder 1.37 \AA . The molecule as a whole is not planar; the nitro groups at C(2), C(3) and C(4) are found to be rotated by 45° , 64° and 19° respectively out of the best plane through the benzene ring.

Introduction

It has been observed that poly-substituted aromatic compounds exhibit distortion of the aromatic nucleus, unusual exocyclic bond lengths, inter- and intramolecular hydrogen bonding, and crowding of substituent groups. Cady & Larson (1965) have reported that 1,3,5-triamino-2,4,6-trinitrobenzene exhibits carbon–carbon bond lengths that average nearly 1.45 \AA . Holden (1962) has reported for 1,3-diamino-2,4,6-trinitrobenzene carbon–nitrogen bond lengths of 1.40 \AA at C(2) and 1.48 \AA at C(6). Both of these molecules form strong hydrogen bond networks, 1,3,5-triamino-2,4,6-trinitrobenzene bonding into sheets and 1,3-diamino-2,4,6-trinitrobenzene into chains with only weak interactions between chains. Gafner & Herbstein (1962) and Ferguson & Sim (1962), among others, have published several papers concerning the effects of overcrowding on aromatic systems.

The structure of 2,3,4,6-tetranitroaniline (TENA) was undertaken to study these phenomena further because of the many possibilities offered by the functional groups present. Vicinal nitro groups are large enough to cause overcrowding and the system as a whole has the possibility of a large degree of conjugation. Because amine groups may act as donors and nitro groups as receptors, the possibilities of inter- and intramolecular hydrogen bonding also exist. Further, since no two nitro groups have the same pair of nearest neighbors, the molecule provides an opportunity for study of the effect of various environments on the configuration of a nitro group. Finally, Flürsheim (1921) has reported enhanced reactivity of the nitro group at C(3) toward certain nucleophiles and it was hoped that this reactivity might be reflected in some feature of the structure.

Determination of the structure

Experimental

A sample of the material was obtained from the Organic Chemistry Division of the U.S. Naval Ordnance Laboratory and suitable crystals were grown without purification by slowly cooling a benzene-acetone solution. The crystal used in the collection of the data was a triangular plate broken from a larger, lath-shaped crystal. This crystal fragment was approximately 0.3 mm in its largest dimension. The linear absorption coefficient for TENA is 15.92 cm^{-1} for $Cu K\alpha$ radiation. No absorption or extinction corrections were applied. The intensity data were collected at room temperature with the crystal rotating about the b axis upon a non-integrating Weissenberg camera using $Cu K\alpha$ radiation ($\lambda = 1.5418 \text{ \AA}$). The reflection intensities were visually estimated using a gray scale prepared from a typical reflection from the same crystal that was used to collect the intensity data. The range of measured intensities was 7 to 41,000 on an arbitrary scale. The number of reflections in the region of reciprocal space measured was 1690 of which 1074 were of detectable intensity.

Crystallographic data

Precession, Weissenberg and oscillation films were used to obtain the crystallographic data for TENA. From these techniques the following data were obtained:

Cell dimensions:

$$a_0 = 7.27 \pm 0.01, b_0 = 11.06 \pm 0.02, c_0 = 12.27 \pm 0.02 \text{ \AA}$$

$$\beta = 98^\circ 8' \pm 0.3'$$

Cell volume: 972 \AA^3 .

Density: observed, 1.87 g.cm^{-3} ; calculated, 1.86 g.cm^{-3} with $Z = 4$. Systematic absences: $h0l$; $l = 2n + 1$; $0k0$; $k = 2n + 1$.

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From these periodic absences on the precession and Weissenberg films the space group was uniquely determined as $P2_1/c$. $F(000)$ is 552.

Trial structure

The visual estimates of intensity were averaged, greater weight being given to those measurements in the mid-region of the intensity range, and $1/L_p$ corrections were applied with the use of a crystallographic computing system, X-ray 63 (Stewart *et al.*, 1964). After trying unsuccessfully to obtain a trial structure through interpretation of the Patterson function we turned to the symbolic addition procedure (Karle & Karle, 1963), which had been successfully applied to several centrosymmetric structures. The first step in this procedure was the conversion of the intensity data to a set of quasi-normalized structure factors, E values (Hauptman & Karle, 1953). The method of conversion was that developed for the X-ray 63 computing system. The atomic scattering factors used in this and all other calculations were those tabulated by Berghuis, Haanappel, Potters, Loopstra, MacGillavry & Veenendaal (1955) for carbon, nitrogen and oxygen. The atomic scattering factors used for hydrogen were those of McWeeny (1951).

E values may be defined as follows:

$$E(hkl) = K F(hkl) \exp [B(\sin \theta(hkl)/\lambda)^x] / \left[\epsilon \sum_{j=1}^N f^2(j) n(j) \right]^{\frac{1}{2}}$$

where

K = scale constant

$F(hkl)$ = relative structure factor

B = overall temperature factor

$\theta(hkl)$ = Bragg angle

λ = X-ray wavelength

x = temperature factor exponent (assumed = 2 for TENA)

ϵ = systematic extinction compensation weight

$f(j)$ = form factor of atom j

$n(j)$ = number of atoms of type j in the cell

N = number of atom types in the cell.

Values of K and B are determined by the conversion program. In this calculation, for each value of B the value of K is set so that the average value of E^2 is 1.0 and the value of B is adjusted by trial and error so as to minimize the resulting sum of $(E^2 - 1)^2$. This has the desired effect of producing a set of quasi-normalized structure factors whose magnitudes tend to show no trend with increasing diffraction angle. These values then constitute a proper set of E values.

The principal advantage of estimating E values by this method is that the data need not be divided into arbitrary groups, but rather may be treated *in toto*. This method, in addition to giving good E values, is a very direct way to estimate temperature and scale factors. The E values are used in the symbolic addition procedure, while $\epsilon(E)^2$ and $\epsilon(E)^2 - 1$ serve as coefficients

for producing sharpened and sharpened, origin-removed Patterson functions.

The symbolic addition procedure was carried out with only those reflections with E values greater than 1.5. An origin was defined by assigning the 612, the 255 and the 368 reflections a phase of 0°, and then using the Σ_2 relationship to assign phases to related reflections. In order to obtain enough phase relationships, it was found necessary to assign the 133, the 240 and the 645 reflections phase cosines of a , b and c respectively. Because of relationships that developed during the application of the symbolic addition procedure it was possible to assign b a value of -1 and c a value of +1. The symbol a was not determined and both the +1 and -1 values were used in subsequently calculated E maps. Of the 114 phases assigned consistent with an a value of +1, 5 were incorrect as shown by the final structure.

E maps, Fourier maps calculated with the quasi-normalized structure factors as coefficients, were calculated for both possible a phase values using these 114 reflections. The 19 largest peaks were selected from each map and refinement begun on both structures.

Preliminary refinement

During the early stages of refinement Fourier methods were employed. Nitrogen atoms, as a compromise between carbon, nitrogen and oxygen atoms, were placed at each indicated atomic position, the electron density calculated at 27 points about each position and the atom moved to coincide with the interpolated position of maximum electron density. Both possible structures were refined by several cycles of this procedure. One structure refined from an R of 0.59 to an R of 0.53, and the other structure from an R of 0.52 to an R of 0.37. A difference Fourier synthesis calculated for this second structure clearly indicated which atoms were carbon, which nitrogen and which oxygen. The atoms were placed as indicated and Fourier refinement continued, several cycles reducing R to 0.19.

Final refinement

Several cycles of full-matrix least squares were run, using individual isotropic temperature factors and level by level scale factors, minimizing $\sum w(\Delta F^2)$. This isotropic refinement reduced R from 0.19 to 0.13. Several more cycles were run with individual anisotropic temperature factors for all non-hydrogen atoms and a Hughes type weighting scheme (Hughes, 1941). The 122, 221, 133, 042 and 240 reflections were given zero weight. The remaining reflections were divided into two classes for weighting purposes. The observed reflections were weighted as follows:

$$w = 10/F_{\text{rel}} \quad F_{\text{rel}} \geq 10,$$

$$w = 1.0 \quad F_{\text{rel}} < 10.$$

Unobserved reflections were weighted

$$w = 1.0 \quad F_c > F_{\text{min}},$$

$$w = 0.0 \quad F_c \leq F_{\text{min}},$$

Table 1. Final atomic parameters

	<i>x</i>	<i>y</i>	<i>z</i>	<i>B</i> ₁₁	<i>B</i> ₂₂	<i>B</i> ₃₃	<i>B</i> ₁₂	<i>B</i> ₁₃	<i>B</i> ₂₃
C(1)	0.7532 (5)	0.3357 (3)	0.2208 (3)	3.0 (1)	2.7 (2)	4.3 (2)	-0.3 (1)	0.3 (1)	0.2 (1)
C(2)	0.5893 (5)	0.4094 (4)	0.2004 (3)	3.7 (2)	2.0 (2)	4.3 (2)	-0.1 (1)	0.8 (1)	0.1 (1)
C(3)	0.4453 (5)	0.3834 (4)	0.1192 (3)	2.6 (1)	3.2 (2)	4.0 (2)	-0.1 (1)	0.6 (1)	0.4 (1)
C(4)	0.4541 (5)	0.2841 (4)	0.0515 (3)	3.6 (2)	3.1 (2)	4.0 (2)	-0.4 (1)	0.2 (1)	0.1 (1)
C(5)	0.6127 (6)	0.2143 (4)	0.0623 (3)	4.4 (1)	2.8 (2)	3.9 (2)	-0.1 (1)	0.5 (1)	0.2 (1)
C(6)	0.7587 (5)	0.2401 (3)	0.1431 (3)	3.4 (2)	2.4 (2)	4.2 (2)	0.1 (1)	0.6 (1)	-0.1 (1)
N(1)	0.8867 (5)	0.3577 (4)	0.3031 (3)	3.9 (1)	3.8 (2)	4.9 (2)	0.6 (1)	-0.5 (1)	-0.9 (1)
N(2)	0.5837 (5)	0.5209 (3)	0.2646 (3)	3.6 (1)	2.7 (2)	5.1 (2)	0.1 (1)	1.0 (1)	-0.1 (1)
N(3)	0.2815 (5)	0.4655 (3)	0.1039 (3)	3.5 (1)	3.6 (2)	5.0 (2)	0.3 (1)	0.8 (1)	0.8 (1)
N(4)	0.2953 (5)	0.2434 (4)	-0.0254 (3)	4.6 (2)	3.9 (2)	4.6 (2)	-1.0 (1)	-0.4 (1)	0.2 (1)
N(6)	0.9187 (5)	0.1591 (3)	0.1489 (3)	4.1 (2)	2.7 (2)	5.5 (2)	0.2 (1)	0.9 (1)	0.0 (1)
O(21)	0.6364 (4)	0.5170 (3)	0.3636 (2)	5.2 (1)	3.8 (2)	4.6 (2)	0.3 (1)	0.5 (1)	-0.8 (1)
O(22)	0.5262 (5)	0.6118 (3)	0.2146 (3)	6.5 (2)	2.7 (2)	6.6 (2)	0.8 (1)	1.2 (1)	0.1 (1)
O(31)	0.1903 (4)	0.4682 (4)	0.1786 (3)	4.7 (2)	7.1 (2)	6.7 (2)	1.6 (1)	2.4 (1)	1.3 (2)
O(32)	0.2551 (5)	0.5215 (3)	0.0185 (3)	4.8 (1)	5.2 (2)	6.2 (2)	1.0 (1)	0.5 (1)	2.2 (1)
O(41)	0.1419 (4)	0.2826 (4)	-0.0150 (3)	3.7 (1)	6.8 (2)	7.6 (2)	-0.6 (1)	-0.6 (1)	-0.2 (2)
O(42)	0.3225 (5)	0.1718 (4)	-0.0977 (3)	6.6 (2)	5.6 (2)	5.5 (2)	-0.7 (1)	-0.8 (1)	-1.4 (1)
O(61)	0.9133 (5)	0.0794 (3)	0.0795 (3)	6.0 (2)	3.2 (2)	6.7 (2)	0.9 (1)	0.5 (1)	-1.4 (1)
O(62)	1.0504 (4)	0.1752 (3)	0.2211 (3)	4.3 (1)	4.8 (2)	6.8 (2)	1.3 (1)	-0.4 (1)	-0.9 (1)
H(5)	0.6125 (72)	0.1431 (51)	0.0098 (41)	4.3*					
H(11)	0.9919 (77)	0.3028 (56)	0.3157 (43)	4.3*					
H(12)	0.8721 (74)	0.4194 (59)	0.3555 (45)	4.3*					

* Isotropic temperature factor used for hydrogen atoms and never refined.

where F_{\min} is F_{rel} calculated from the minimum observed intensity. The final parameters are shown in Table 1. The anisotropic temperature factors are defined by the following equation:

$$\begin{aligned} \text{T.F.} = & \exp [-\frac{1}{4}(h^2a^{*2}B_{11} + k^2b^{*2}B_{22} + l^2c^{*2}B_{33} \\ & + 2hka^*b^*B_{12} + 2hla^*c^*B_{13} + 2klb^*c^*B_{23})] . \end{aligned}$$

These B_{ij} values are on the same scale as the overall isotropic B in the expression:

$$F_o = F_e \exp (-B \sin^2 \theta / \lambda^2) .$$

The errors, which are in the last place or last two places, are given in parentheses. The agreement factors from the last least-squares cycle are found in Table 2, and the structure factors in Table 3.

Table 2. Values from final least-squares cycle

$\Sigma w(\Delta F)^2$	665
$[\Sigma w(\Delta F)^2]/(\text{NO}-\text{NV})^+$	0.66
R , omitting unobserveds	0.081
R , with reflection multiplicity omitting unobserveds	0.081
Weighted R value, omitting unobserveds	0.104
R value of unobserveds calculating greater than intensity minimum	0.057
Number of unobserveds calculating greater than intensity minimum	84 of 616
Total number of observations (NO)	1690
Total number of variables (NV)	190

Hydrogen atoms

The hydrogen atoms were located from difference syntheses and their positions refined by full-matrix least-squares along with the other parameters. The hydrogen atoms were assigned an isotropic temperature

factor which was not refined. The value chosen was that indicated for the overall B by the calculation of quasi-normalized structure factors. The atomic positions of the hydrogen atoms were quite stable, no damping or resetting being necessary.

Scale factors

Level scale factors, where each level is that group of reflections with constant index k , were used throughout. No cross reference exposures were used, but an effort was made to keep the data on the same scale by exposing all of the film packs for the same length of time, at the same current and voltage and developing them in the same manner. A first approximation of the value of an overall scale factor was obtained from the calculation of quasi-normalized structure factors by the method described above.

The overall scale factor was used during Fourier refinement and as the starting value for the least-squares refinement. The values of the level scale factors were refined only during the first cycles of least-squares refinement along with individual isotropic atomic temperature factors. No further refinement of the scale factors was attempted during refinement with anisotropic temperature factors because of the high correlation between the scale factors and the B_{22} temperature factors.

The final level scale factors obtained in this way show no trend with k and the values of the B_{22} are not noticeably different from the B_{11} and B_{33} values.

Lingafelter & Donohue (1966) have pointed out that refinement of level scale factors along one crystallographic direction is undesirable because it increases the standard deviation of the positional and thermal parameters associated with that direction.

than O(31) or O(32) (Fig. 1). These distances are N(3) ··· O(41), 2.61 Å; O(32) ··· O(41), 2.78 Å; N(3) ··· O(22), 2.62 Å; and O(31) ··· O(22), 2.87 Å. The angles defined by O-N-O are found to be approximately 124° with the exception of the angle defined by O(31)-N(3)-O(32), which is 128°. This larger angle is possible because the nitro group is rotated so that its oxygen atoms are not close to neighboring atoms.

Table 4. Distance of atoms in Fig. 1 from plane of the six carbon atoms

C(1)	-0.04 Å	O(61)	0.08 Å
C(2)	0.02	O(62)	0.01
C(3)	0.01	H(5)	-0.01
C(4)	-0.03	H(11)	-0.18
C(5)	0.01	H(12)	-0.20
C(6)	0.02	N(4')	-0.72
N(1)	-0.10	O(41')	-1.64
N(2)	0.18	O(42')	0.46
N(3)	0.08	N(6'')	2.18
N(4)	-0.22	O(61'')	1.08
N(6)	0.03	O(62'')	2.41
O(21)	-0.50	N(1''')	-1.63
O(22)	0.99	H(11''')	-2.27
O(31)	-0.88	H(12''')	-0.73
O(32)	1.08	N(1 ^{iv})	-0.92
O(41)	-0.01	H(11 ^{iv})	-0.28
O(42)	-0.10	H(12 ^{iv})	-1.83

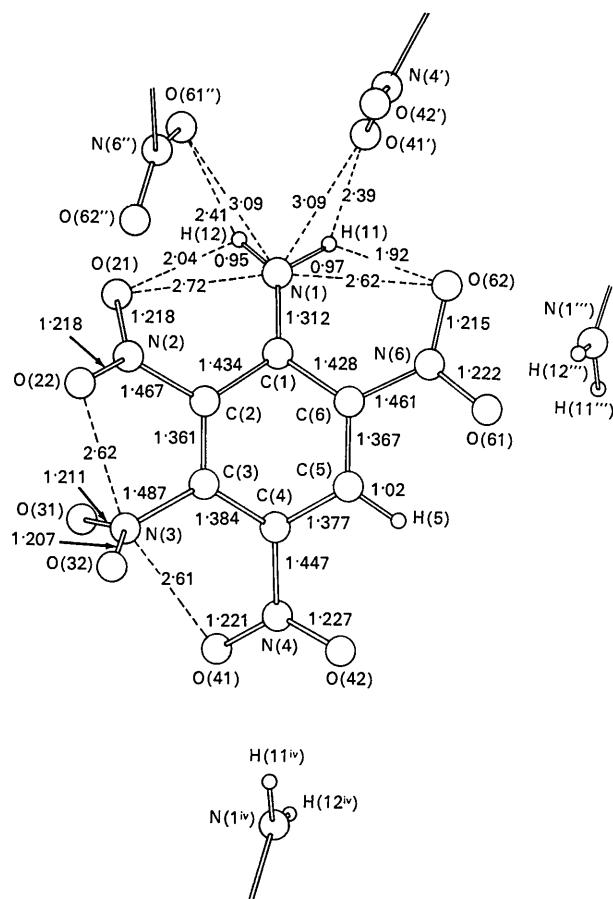


Fig. 1. Bond lengths in tetranitroaniline.

The non-bonded distances of interest at the amine group are O(61) ··· N(1), 2.62 Å and O(21) ··· N(1), 2.72 Å. Neither of these distances appear to be unusually short when compared with other non-bonded N ··· O distances in the molecule. It should also be noted that in 1,3,5-triamino-2,4,6-trinitrobenzene and 1,3-diamino-2,4,6-trinitrobenzene the N ··· O distances considered to indicate intramolecular hydrogen-bonding are more nearly 2.50 than 2.60 Å. These observations indicate that if there is intramolecular hydrogen bonding in TENA, it is of moderate strength. However, even a weak hydrogen bond of the amino group to the nitro group at C(6) might tend to keep that nitro group in a more coplanar configuration with the ring.

Intermolecular distances less than the sum of the van der Waals radii given by Pauling (1960) include H(11) to O(42') of the molecule at $(1+x, \frac{1}{2}-y, \frac{1}{2}+z)$, 2.49 Å; O(21) to C(5^{vi}) of the molecule at $(1-x, \frac{1}{2}+y, \frac{1}{2}-z)$, 3.06 Å; and O(21) to H(5^{vii}) of the molecule at $(x, \frac{1}{2}-y, \frac{1}{2}+z)$, 2.54 Å. Oxygen atom O(32^{viii}) of the molecule at $(1-x, 1-y, -z)$ is approximately centered over the ring. The largest separation between this atom and a carbon atom of the ring is C(1) ··· O(32^{viii}), 3.33 Å, while the shortest distance is C(3) ··· O(32^{viii}), 3.14 Å.

The distances between N(1) and O(41') of the molecule at $(1+x, \frac{1}{2}-y, \frac{1}{2}+z)$ and O(61'') of the molecule at $(2-x, \frac{1}{2}+y, \frac{1}{2}-z)$ are 3.09 Å in each case. The corresponding H ··· O distances are 2.36 and 2.42 Å. This bonding corresponds very closely to that in *p*-nitroaniline where the amino group is hydrogen bonded to nitro groups of two different molecules. In *p*-nitroaniline the N ··· O distances are 3.07 and 3.14 Å. These two intermolecular hydrogen-bonding networks differ in that in *p*-nitroaniline both oxygen atoms of the nitro group are involved in intermolecular hydrogen bonds to the amine groups of separate molecules, whereas in TENA the nitro groups at C(4) and C(6) have only one oxygen atom each involved in an intermolecular hydrogen bond to the amino groups of different molecules.

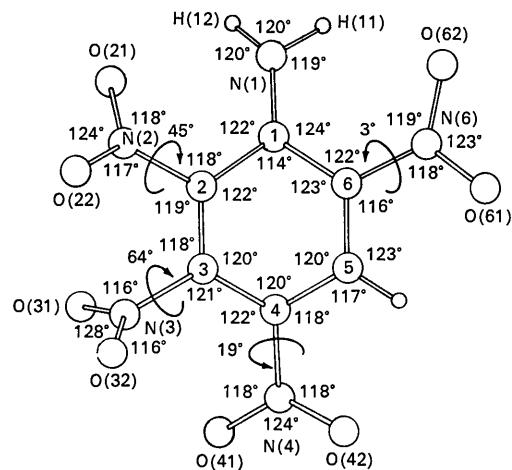


Fig. 2. Bond angles in tetranitroaniline.

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The Crystal Structures of UC_2^*

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A high-temperature neutron diffraction study of UC_2 shows a tetragonal phase at 1700°C with the $C11\alpha$ calcium carbide type structure, $a_0=3.633$ Å, $c_0=6.036$ Å, $z=0.395$, and a cubic phase at 1900°C with the $B1$ sodium chloride type structure, $a_0=5.488$ Å. The cubic phase contains C_2 groups whose alignment was reduced to two possible cases, free rotation or random orientation along [111].

The phase diagram of the uranium–carbon system (Storms, 1965) lists a compound, uranium dicarbide, of approximate composition UC_2 existing in two high-temperature forms but unstable at room temperature. A structure has been described from room temperature studies of quenched samples (Litz, Garrett & Croxton, 1948; Austin, 1959) as tetragonal, space group $I4/mmm$, with uranium atoms in $(0, 0, 0)$, carbon atoms in $(0, 0, z)$, $z=0.388$. The structure of the high temperature form has been reported from high temperature X-ray data to be the CaF_2 type (Wilson, 1960) or the FeS_2 type (Bredig, 1960).

A sample of uranium dicarbide containing excess carbon was prepared from the high-purity elements by arc-melting and pulverizing, and was examined by high-temperature neutron-diffraction (Bowman, Witteman, Arnold, Hull & Bowman, 1966). The diffraction data were obtained from $2\theta=25$ to 70° at a wavelength of 1.3926 Å, and were corrected by subtracting a blank run on the graphite sample holder. The tetra-

gonal form was observed at 1700° and the cubic form at 1900°C . The corrected patterns are shown in Fig. 1. Observed intensities were determined by least-squares analysis of the diffraction data (Bowman, Wallace, Yarnell, Wenzel & Storms, 1965), and were fitted to trial structures by least-squares solutions of the equation (Bacon, 1962)

$$I = K \frac{\exp(-\mu t \sec \theta)}{\sin^2 2\theta} \exp\left(-2B \frac{\sin^2 \theta}{\lambda^2}\right) jF^2. \quad (1)$$

Chemical analysis of the sample showed U 90.2%, C 9.40%, free C 0.96%, O 0.15%, N 60 ppm. Accordingly, a composition of $\text{UC}_{1.85}$ was used for the intensity calculations, with scattering factors of 0.85×10^{-12} cm for uranium and 0.661×10^{-12} cm for carbon.

The tetragonal structure previously reported was confirmed, with $K=0.089 \pm 0.007$, $B=2.7 \pm 0.6$, $z=0.395 \pm 0.003$, $R=0.07$, where $R=\sum w|I_o - I_{cl}|/\sum wI_o$. Lattice parameters at 1700° were $a_0=3.633$, $c_0=6.036$ Å, giving a C–C distance of 1.27 ± 0.04 Å.

The high-temperature pattern was indexed as cubic, $a_0=5.488$ Å at 1900° , in agreement with the earlier X-ray data (Wilson, 1960). The data require a sodium

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